

Fig. 8. Comparison of present and previous high-temperature K versus T results for Ge.

of the two experiments. The reason for this probably real discrepancy is not completely understood.

INTERPRETATION OF LOW-TEMPERATURE MEASUREMENTS

In the temperature range below 300°K, the K results on Si and Ge can be interpreted using Callaway's formalism.³⁷⁻³⁹ The K due to heat transport by phonons is

TABLE II. Literature references to the thermal conductivity of Si and Ge for $T \geq 300^\circ\text{K}$.

Mtrl	Author	Method	Temp. range °K	Year	Ref.
Si	Kohnigsberger and Wels	c	293	1911	14
	Kuprovski and Geld	b	380-1190	1956	15, 16
	Kingery	c	370-540	1959	17
	Mette <i>et al.</i>	e	550-770	1960	18
	Stuckes	c	300-580	1960	19
	Morris and Huest	c	300-700	1961	20
	Abeles <i>et al.</i>	d	310-1200	1962	21, 22
	Shanks <i>et al.</i>	d	300-1400	1963	23
	Morris and Martin	c	680-1000	1963	24
	Ge	Greico and Montgomery	c	298	1952
McCarthy and Ballard		c	280-370	1955	27
Ioffe		a	300-870	1956	28
Shtenbeck and Baranskii		a	200-370	1957	29
Baranskii and Konoplasova		a	80-370	1958	30
Abeles		a	320-1070	1959	31
Kettel		a	390-1000	1959	32
Pankove		e	370-970	1959	33
Stuckes		c	310-680	1960	34
Devyatkova and Smirnov		a	80-440	1960, 62	35, 36
Slack and Glassbrenner		a, b	3-1020	1960	6
Abeles <i>et al.</i>		d	310-1070	1962	21, 22

- ^a Absolute, steady-state, longitudinal heat flow.
- ^b Absolute, steady-state, radial heat flow.
- ^c Comparative, steady-state longitudinal heat flow.
- ^d Variable state, thermal diffusivity.
- ^e Other.

³⁷ J. Callaway, Phys. Rev. 113, 1046 (1959).
³⁸ J. Callaway and H. C. von Baeyer, Phys. Rev. 120, 1149 (1960).
³⁹ J. Callaway, Phys. Rev. 122, 787 (1961).

given by

$$K = \frac{k}{2\pi^2 v} \left(\frac{kT}{\hbar} \right)^3 \int_0^{\theta/T} \tau_C(x, T) \frac{x^4 e^x}{(e^x - 1)^2} dx, \quad (1)$$

where k = Boltzmann's constant, v =average sound velocity, T =absolute temperature, \hbar =Planck's constant, θ =Debye temperature, $x = (\hbar\omega/kT)$, ω =phonon frequency, and τ_C is a combined relaxation time. The complicating effects³⁷ of normal phonon scattering have been neglected. The combined relaxation time is thus taken as

$$\tau_C^{-1} = \tau_B^{-1} + \tau_I^{-1} + \tau_U^{-1}, \quad (2)$$

where τ_B , τ_I , τ_U are the relaxation times for boundary, isotope, and umklapp scattering of the phonons.

The particular form of the relaxation times shall now be examined. Klemens^{40,41} has derived a relaxation time for scattering by point imperfections. His results are valid for isotope scattering.

$$\tau_I^{-1}(\omega) = 3V_0 \Gamma \omega^4 / \pi v^3, \quad (3)$$

where V_0 =average volume per atom in the crystal, and Γ =point impurity scattering parameter.⁴² For umklapp scattering the following semiempirical expression⁴³ will be used:

$$\tau_U^{-1} = B_U \omega^2 T. \quad (4)$$

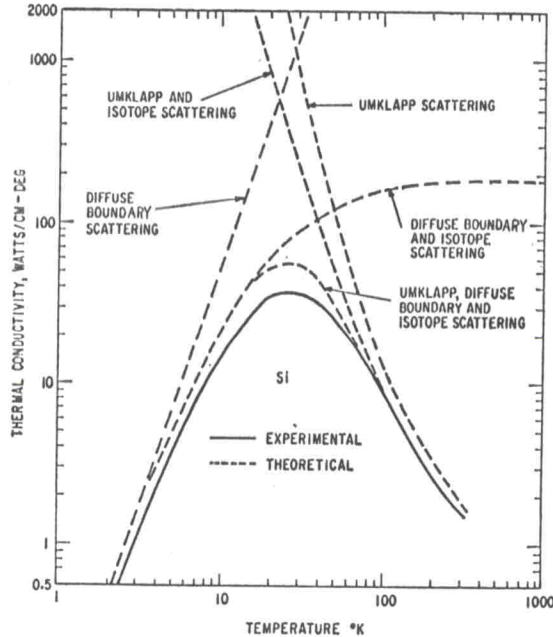


Fig. 9. Comparison of the low-temperature K results for Si with the theory in which various phonon scattering mechanisms are considered.

⁴⁰ P. G. Klemens, in *Solid State Physics*, edited by F. Seitz and D. Turnbull (Academic Press Inc., New York, 1958), Vol. 7, p. 1.

⁴¹ P. G. Klemens, in *Handbuch der Physik*, edited by S. Flügge (Springer-Verlag, Berlin, 1956), Vol. 14, p. 198.

⁴² G. A. Slack, Phys. Rev. 126, 427 (1962).

⁴³ G. A. Slack and S. Galginaitis, Phys. Rev. 133, A253 (1964).